

On Some Nonlocal in Time and Space Parabolic Problem

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Abstract

The goal of this note is to study nonlinear parabolic problems nonlocal in time and space. We first establish the existence of a solution and its uniqueness in certain cases. Finally we consider its asymptotic behaviour.

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1 Introduction and notation

We will denote by Ω a smooth bounded open set of \mathbb{R}^n , $n \geq 2$ with boundary $\partial\Omega$. We would like to consider the following problem. Find $u = u(x, t)$ such that

$$\begin{cases} u_t - \alpha \left(\int_{\Omega} g(x) u(x, t) dx \right) \Delta u + \beta \left(\int_0^t h(s) u(x, s) ds \right) u = f & \text{in } \Omega \times (0, T), \\ u(\cdot, t) = 0 \text{ on } \partial\Omega, \quad t \in (0, T), \quad u(x, 0) = u_0(x). \end{cases} \quad (1.1)$$

T is a positive number, u_0, f, g, h are given data. The equation could be regarded as a model of population dynamics where $u(x, t)$ is the density of a population at the location x , at the time t . The nonlinear terms are an account for a death or diffusion rate which at time t depends on the total population having been at the location x in the past or in Ω at time t . To study this issue we were inspired by the papers [7]-[10] where a similar problem was introduced at the difference that in (1.1) the integral goes up to T which we think is an interesting point of view but perhaps a bit surprising from a realistic one in our framework.

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The paper is divided as follows. In the next section we prove existence of a weak solution to (1.1). In the subsequent section we establish a result of uniqueness. Note that in comparison to [7]-[10] our result is global. Finally we study in a simple case the asymptotic behaviour of the solution to (1.1).

2 A result of existence

We denote by $L^p(\Omega)$, $1 \leq p \leq +\infty$ the usual L^p -space on Ω . It is equipped with its usual norm and for instance, in the case where $p = 2$, we denote it by $\|\cdot\|_2$ i.e.

$$\|v\|_2^2 = \int_{\Omega} v(x)^2 dx \quad \forall v \in L^2(\Omega).$$

We refer the reader to [2]-[6] for the notation used in the sequel, for instance for $H_0^1(\Omega)$ or its dual $H^{-1}(\Omega)$ or the spaces $L^2(0, T; V)$, $L^2(0, T; V')$ when V is a Banach space.

The main result of this section is the following :

Theorem 2.1. *Set $V = H_0^1(\Omega)$, $V' = H^{-1}(\Omega)$. Suppose*

$$u_0 \in L^2(\Omega), f \in L^2(0, T; V'), \alpha, \beta \in C(\mathbb{R}) \cap L^\infty(\mathbb{R}), g \in L^\infty(\Omega), h \in L^\infty(0, T), \forall T, \quad (2.1)$$

and that for some positive constant a one has

$$0 < a \leq \alpha. \quad (2.2)$$

Then there exists a weak solution to (1.1). $C(\mathbb{R})$ denotes the space of continuous functions.

Proof. 1. One can assume that $\beta \geq 1$.

Indeed, suppose that we can solve (1.1) in this case. u is solution to (1.1) iff

$$\tilde{u} = e^{-\lambda t} u \quad (2.3)$$

satisfies

$$\begin{aligned} & (e^{\lambda t} \tilde{u})_t - \alpha \left(\int_{\Omega} g(x) e^{\lambda t} \tilde{u}(x, t) dx \right) e^{\lambda t} \Delta \tilde{u} + \beta \left(\int_0^t h(s) e^{\lambda s} \tilde{u}(x, s) ds \right) e^{\lambda t} \tilde{u} = f, \\ \Leftrightarrow & e^{\lambda t} \tilde{u}_t + \lambda e^{\lambda t} \tilde{u} - \alpha \left(\int_{\Omega} g(x) e^{\lambda t} \tilde{u}(x, t) dx \right) e^{\lambda t} \Delta \tilde{u} + \beta \left(\int_0^t h(s) e^{\lambda s} \tilde{u}(x, s) ds \right) e^{\lambda t} \tilde{u} = f, \\ \Leftrightarrow & \tilde{u}_t - \alpha \left(\int_{\Omega} g(x) e^{\lambda t} \tilde{u}(x, t) dx \right) \Delta \tilde{u} + \left\{ \lambda + \beta \left(\int_0^t h(s) e^{\lambda s} \tilde{u}(x, s) ds \right) \right\} \tilde{u} = e^{-\lambda t} f, \end{aligned}$$

i.e. iff \tilde{u} satisfies (1.1) with f, h, α replaced respectively by $e^{-\lambda t} f, e^{\lambda t} h, \alpha(e^{\lambda t} \cdot)$ and β by $\lambda + \beta$ which is greater than 1 for λ large enough.

2. We suppose that $\beta \geq 1$.

Let $w \in L^2(0, T; L^2(\Omega)) \subset L^1(0, T; L^1(\Omega))$. Then, see [2], there exists a unique $u = S(w)$ solution to

$$\begin{cases} u \in L^2(0, T; V), \quad u_t \in L^2(0, T; V'), \\ \frac{d}{dt}(u, v) + \alpha \left(\int_{\Omega} g(x)w(x, t)dx \right) \int_{\Omega} \nabla u \cdot \nabla v dx + \left(\beta \left(\int_0^t h(s)w(x, s)ds \right) u, v \right) \\ = \langle f, v \rangle \quad \forall v \in H_0^1(\Omega), \text{ in } \mathcal{D}'(0, T). \end{cases} \quad (2.4)$$

In the equation above we denote by (\cdot, \cdot) the canonical scalar product in $L^2(\Omega)$ and by $\langle \cdot, \cdot \rangle$ the duality between $H^{-1}(\Omega)$ and $H_0^1(\Omega)$, $\mathcal{D}(0, T)$ and $\mathcal{D}'(0, T)$ denote respectively the space of C^∞ functions with compact support in $(0, T)$ and its dual, the usual space of distributions on $(0, T)$. (Cf. for instance [2]). We will be done if we can show that S has a fixed point. Taking in the equation above $v = u$ we get easily if $a \wedge 1$ denotes the minimum of a and 1

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} |u|_2^2 + a \wedge 1 \int_{\Omega} (|\nabla u|^2 + u^2) dx &\leq \langle f, u \rangle \leq |f|_{V'} \|\nabla u\|_2 \\ &\leq \frac{1}{2(a \wedge 1)} |f|_{V'}^2 + \frac{a \wedge 1}{2} \|\nabla u\|_2^2. \end{aligned}$$

$|f|_{V'}$ denotes the strong dual norm of f in $H^{-1}(\Omega)$ associated to the norm $\|\nabla u\|_2$ in $H_0^1(\Omega)$. From this we derive

$$\frac{d}{dt} |u|_2^2 + (a \wedge 1) \int_{\Omega} (|\nabla u|^2 + u^2) dx \leq \frac{1}{a \wedge 1} |f|_{V'}^2$$

and after an integration in t

$$|u|_2^2 + (a \wedge 1) \int_0^t \int_{\Omega} (|\nabla u|^2 + u^2) dx ds \leq |u_0|_2^2 + \frac{1}{a \wedge 1} \int_0^t |f(\cdot, s)|_{V'}^2 ds.$$

It follows that

$$|u|_{L^2(0, T; V)}, |u|_{L^2(0, T; L^2(\Omega))} \leq C^2 = \frac{1}{(a \wedge 1)} \left(|u_0|_2^2 + \frac{1}{a \wedge 1} \int_0^T |f(\cdot, s)|_{V'}^2 ds \right).$$

Set

$$B = \{v \in L^2(0, T; L^2(\Omega)) \mid |v|_{L^2(0, T; L^2(\Omega))} \leq C\}.$$

Clearly, S maps B into itself. Moreover since

$$u_t = \alpha \left(\int_{\Omega} g(x)w(x, t)dx \right) \Delta u - \beta \left(\int_0^t h(s)w(x, s)ds \right) u + f \text{ in } V'$$

u_t is uniformly bounded in $L^2(0, T; V')$ and $S(B)$ is relatively compact in B . The existence of a weak solution to (1.1) will follow by the Schauder fixed point theorem if S is continuous. To show that, let $w_n \in L^2(0, T; L^2(\Omega))$ such that

$$w_n \rightarrow w \text{ in } L^2(0, T; L^2(\Omega)).$$

Denote $u_n = S(w_n)$. The estimates above hold and one can extract a subsequence such that, if we still label it by n

$$\begin{aligned}
gw_n &\rightarrow gw \text{ in } L^2(0, T; L^2(\Omega)), \\
hw_n &\rightarrow hw \text{ in } L^2(0, T; L^2(\Omega)), \\
u_n &\rightarrow u_\infty \text{ in } L^2(0, T; L^2(\Omega)), \\
\nabla u_n &\rightharpoonup \nabla u_\infty \text{ in } L^2(0, T; L^2(\Omega)), \\
(u_n)_t &\rightharpoonup (u_\infty)_t \text{ in } L^2(0, T; V').
\end{aligned} \tag{2.5}$$

By definition of u_n we have for every $v \in H_0^1(\Omega)$ and every $\varphi \in \mathcal{D}(0, T)$

$$\begin{aligned}
\int_0^T -(u_n, v)\varphi'(t)dt + \int_0^T \varphi(t) \int_\Omega \alpha \left(\int_\Omega g(x)w(x, t)dx \right) \nabla u_n \cdot \nabla v \, dxdt \\
+ \int_0^T \varphi(t) \int_\Omega \beta \left(\int_0^t h(s)w_n(x, s)ds \right) u_n v \, dxdt = \int_0^T \langle f, v \rangle \varphi(t)dt.
\end{aligned}$$

By the Lebesgue theorem

$$\varphi(t)\beta \left(\int_0^t h(s)w_n(x, s)ds \right) v \rightarrow \varphi(t)\beta \left(\int_0^t h(s)w(x, s)ds \right) v \text{ in } L^2(0, T; L^2(\Omega)). \tag{2.6}$$

Indeed, note that

$$\begin{aligned}
\left| \int_0^t h(s)w_n(x, s)ds - \int_0^t h(s)w(x, s)ds \right| \\
\leq \int_0^t |h|_\infty |w_n - w|(x, s)ds \\
\leq |h|_\infty \sqrt{T} \left\{ \int_0^T (w_n - w)^2(x, s)ds \right\}^{\frac{1}{2}} \rightarrow 0 \text{ a.e.}
\end{aligned}$$

up to a subsequence. $|h|_\infty$ is the $L^\infty(0, T)$ -norm of h . Using (2.6) and the analogue written for α and g , one can pass to the limit in the equation satisfied by u_n . It follows that $u_\infty = S(w)$. Since the limit of u_n is unique the whole sequence u_n converges toward $u_\infty = S(w)$ and thus S is continuous. This completes the proof of the theorem. \square

Remark 1. *The same existence result holds if in (1.1) one replaces the integral on $(0, t)$ by*

$$\int_0^{t'} h(s)u(x, s)ds$$

where t' is any real number in $(0, T]$.

3 Uniqueness issue

One has the following estimate for the solution to (1.1):

Proposition 3.1. *Suppose that $u_0 \in L^\infty(\Omega)$, $f \in L^\infty(\Omega \times (0, T))$, $\beta \geq 1$. Then it holds*

$$|u| \leq K = |f|_\infty \vee |u_0|_\infty. \quad (3.1)$$

(\vee stands for the maximum of two numbers).

Proof. One has

$$\frac{d}{dt}(u - K) - \nabla \cdot \left(\alpha \left(\int_\Omega g u \, dx \right) \nabla(u - K) \right) + \beta \left(\int_0^t h u \, ds \right) u - K = f - K \leq 0.$$

It follows, using as test function $(u - K)^+$ where $(\)^+$ denotes the positive part of a function

$$\frac{1}{2} \frac{d}{dt} |(u - K)^+|_2^2 + a \wedge 1 \int_\Omega |\nabla(u - K)^+|^2 + ((u - K)^+)^2 \leq 0.$$

This implies that

$$\frac{d}{dt} (|(u - K)^+|_2^2 e^{2(a \wedge 1)t}) \leq 0$$

and since this quantity vanishes at 0 it vanishes for all time. This shows that $u \leq K$. Since $-u$ satisfies a similar equation one has also $-u \leq K$. This completes the proof of the proposition. \square

One can then prove the following uniqueness result :

Theorem 3.1. *Suppose that $u_0 \in L^\infty(\Omega)$, $f \in L^\infty(\Omega \times (0, T))$, $g \in L^\infty(\Omega)$, $h \in L^\infty(0, T)$. Suppose that $\beta \geq 1$, α are Lipschitz continuous in the sense that for some positive constant C_α , C_β*

$$|\alpha(\xi) - \alpha(\eta)| \leq C_\alpha |\xi - \eta|, \quad |\beta(\xi) - \beta(\eta)| \leq C_\beta |\xi - \eta| \quad \forall \xi, \eta \in \mathbb{R}, \quad (3.2)$$

then the weak solution to (1.1) is unique.

Proof. Let u_1, u_2 be two solutions to (1.1). By subtraction one gets

$$\begin{aligned} \frac{d}{dt}(u_1 - u_2) - \alpha \left(\int_\Omega g(x) u_1(x, t) \, dx \right) \Delta(u_1 - u_2) + \beta \left(\int_0^t h(s) u_1(x, s) \, ds \right) (u_1 - u_2) \\ = \left(\alpha \left(\int_\Omega g(x) u_1(x, t) \, dx \right) - \alpha \left(\int_\Omega g(x) u_2(x, t) \, dx \right) \right) \Delta u_2 \\ - \left\{ \beta \left(\int_0^t h(s) u_1(x, s) \, ds \right) - \beta \left(\int_0^t h(s) u_2(x, s) \, ds \right) \right\} u_2. \end{aligned}$$

Multiplying by $(u_1 - u_2)$ and integrating on Ω we get

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} |u_1 - u_2|_2^2 + \alpha \left(\int_\Omega g(x) u_1(x, t) \, dx \right) \int_\Omega |\nabla(u_1 - u_2)|^2 \, dx \\ + \int_\Omega \beta \left(\int_0^t h(s) u_1(x, s) \, ds \right) (u_1 - u_2)^2 \, dx \\ = - \int_\Omega \left(\alpha \left(\int_\Omega g(x) u_1(x, t) \, dx \right) - \alpha \left(\int_\Omega g(x) u_2(x, t) \, dx \right) \right) \nabla u_2 \cdot \nabla(u_1 - u_2) \, dx \\ - \int_\Omega \left\{ \beta \left(\int_0^t h(s) u_1(x, s) \, ds \right) - \beta \left(\int_0^t h(s) u_2(x, s) \, ds \right) \right\} u_2 (u_1 - u_2) \, dx. \end{aligned}$$

By (2.2) and since $\beta \geq 1$ we derive

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} |u_1 - u_2|_2^2 + a \int_{\Omega} |\nabla(u_1 - u_2)|^2 dx + \int_{\Omega} (u_1 - u_2)^2 dx \\ & \leq \int_{\Omega} |\alpha \left(\int_{\Omega} g(x) u_1(x, t) dx \right) - \alpha \left(\int_{\Omega} g(x) u_2(x, t) dx \right)| |\nabla u_2| |\nabla(u_1 - u_2)| dx \\ & \quad + \int_{\Omega} |\{\beta \left(\int_0^t h(s) u_1(x, s) ds \right) - \beta \left(\int_0^t h(s) u_2(x, s) ds \right)\}| |u_2| |(u_1 - u_2)| dx, \end{aligned}$$

which implies

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} |u_1 - u_2|_2^2 + a \int_{\Omega} |\nabla(u_1 - u_2)|^2 dx + \int_{\Omega} (u_1 - u_2)^2 dx \\ & \leq \int_{\Omega} C_{\alpha} \left| \int_{\Omega} g(x) (u_1(x, t) - u_2(x, t)) dx \right| |\nabla u_2| |\nabla(u_1 - u_2)| dx \\ & \quad + \int_{\Omega} C_{\beta} \left| \int_0^t h(s) (u_1(x, s) - u_2(x, s)) ds \right| |u_2| |(u_1 - u_2)| dx \\ & \leq \int_{\Omega} C_{\alpha} g_{\infty} \left(\int_{\Omega} |u_1(x, t) - u_2(x, t)| dx \right) |\nabla u_2| |\nabla(u_1 - u_2)| dx \\ & \quad + \int_{\Omega} C_{\beta} h_{\infty} \int_0^t |u_1(x, s) - u_2(x, s)| ds |u_2(x, t)| |(u_1 - u_2)(x, t)| dx \end{aligned}$$

where g_{∞} , h_{∞} denote the $L^{\infty}(\Omega)$ and $L^{\infty}(0, T)$ norms of g and h . Now we use (3.1) and the Young inequality

$$ab \leq \epsilon a^2 + C_{\epsilon} b^2$$

to get

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} |u_1 - u_2|_2^2 + a \int_{\Omega} |\nabla(u_1 - u_2)|^2 dx + \int_{\Omega} (u_1 - u_2)^2 dx \\ & \leq \int_{\Omega} C_{\epsilon} \left\{ C_{\alpha} g_{\infty} \left(\int_{\Omega} |u_1(x, t) - u_2(x, t)| dx \right) |\nabla u_2| \right\}^2 + \epsilon |\nabla(u_1 - u_2)|^2 dx \\ & \quad + \int_{\Omega} C_{\beta} h_{\infty} K \frac{1}{2} \frac{d}{dt} \left(\int_0^t |u_1(x, s) - u_2(x, s)| ds \right)^2 dx. \end{aligned}$$

Choosing $\epsilon = \frac{a}{2}$ it comes

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} |u_1 - u_2|_2^2 + \frac{a}{2} \int_{\Omega} |\nabla(u_1 - u_2)|^2 dx + \int_{\Omega} (u_1 - u_2)^2 dx \\ & \leq \int_{\Omega} C_{\epsilon} C_{\alpha}^2 g_{\infty}^2 |\nabla u_2|^2 \left(\int_{\Omega} |u_1(x, t) - u_2(x, t)| dx \right)^2 dx \\ & \quad + \int_{\Omega} C_{\beta} h_{\infty} K \frac{1}{2} \frac{d}{dt} \left(\int_0^t |u_1(x, s) - u_2(x, s)| ds \right)^2 dx \\ & \leq \int_{\Omega} C_{\epsilon} C_{\alpha}^2 g_{\infty}^2 |\nabla u_2|^2 |\Omega| |u_1 - u_2|_2^2 dx \\ & \quad + \int_{\Omega} C_{\beta} h_{\infty} K \frac{1}{2} \frac{d}{dt} \left(\int_0^t |u_1(x, s) - u_2(x, s)| ds \right)^2 dx. \end{aligned}$$

We used Hölder's inequality, $|\Omega|$ denotes the measure of Ω . Thus we obtain

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} |u_1 - u_2|_2^2 + \frac{a}{2} \int_{\Omega} |\nabla(u_1 - u_2)|^2 dx + \int_{\Omega} (u_1 - u_2)^2 dx \\ \leq C_{\epsilon} C_{\alpha}^2 g_{\infty}^2 |\nabla u_2|_2^2 |\Omega| |u_1 - u_2|_2^2 \\ + \int_{\Omega} C_{\beta} h_{\infty} K \frac{1}{2} \frac{d}{dt} \left(\int_0^t |u_1(x, s) - u_2(x, s)| ds \right)^2 dx. \end{aligned}$$

Integrating between 0 and t we derive

$$\begin{aligned} |u_1 - u_2|_2^2 &\leq 2 \int_0^t C_{\epsilon} C_{\alpha}^2 g_{\infty}^2 |\nabla u_2|_2^2 |\Omega| |u_1 - u_2|_2^2 ds \\ &\quad + \int_{\Omega} C_{\beta} h_{\infty} K \left(\int_0^t |u_1(x, s) - u_2(x, s)| ds \right)^2 dx \\ &\leq 2 \int_0^t C_{\epsilon} C_{\alpha}^2 g_{\infty}^2 |\nabla u_2|_2^2 |\Omega| |u_1(x, t) - u_2(x, t)|_2^2 dt \\ &\quad + \int_{\Omega} C_{\beta} h_{\infty} K t \int_0^t |u_1(x, s) - u_2(x, s)|^2 ds dx \\ &= \int_0^t \left(2C_{\epsilon} C_{\alpha}^2 g_{\infty}^2 |\Omega| |\nabla u_2|_2^2 + C_{\beta} h_{\infty} K t \right) |u_1 - u_2|_2^2 dt. \end{aligned}$$

Since $\left(2C_{\epsilon} C_{\alpha}^2 g_{\infty}^2 |\Omega| |\nabla u_2|_2^2 + C_{\beta} h_{\infty} K t \right) \in L^1(0, T)$ the result follows from the Gronwall inequality. \square

4 Stationary problem

In this section we consider u solution to (1.1) and we assume

$$f, u_0, g, h \geq 0. \quad (4.1)$$

Moreover we assume that

$$\beta(z) \text{ admits a limit when } z \rightarrow +\infty. \quad (4.2)$$

First notice that (4.1) implies that $u \geq 0$. Indeed multiplying (1.1) by $-u^-$ we get

$$\frac{1}{2} \frac{d}{dt} |u^-|_2^2 + \alpha \left(\int_{\Omega} g u dx \right) \int_{\Omega} |\nabla u^-|^2 dx + \int_{\Omega} \beta \left(\int_0^t h u ds \right) (u^-)^2 dx = -(f, u^-) \leq 0.$$

Since α, β are positive we get

$$\frac{1}{2} \frac{d}{dt} |u^-|_2^2 \leq 0$$

i.e. $u^- = 0$ since $u^-(x, 0) = 0$. Since $u \geq 0$, then

$$t \rightarrow \int_0^t h(s) u(x, s) ds$$

is nondecreasing in time and has a limit when $t \rightarrow +\infty$ for almost every $x \in \Omega$ and so does

$$\beta \left(\int_0^t h(s) u(x, s) ds \right).$$

We denote by $\beta_\infty(x) \in L^\infty(\Omega)$ this limit. Then the stationary problem associated to (1.1) is : find u_∞ weak solution to

$$\begin{cases} -\alpha \left(\int_\Omega g(x) u_\infty(x) dx \right) \Delta u_\infty + \beta_\infty u_\infty = f(x) \text{ in } \Omega, \\ u_\infty = 0 \text{ on } \partial\Omega. \end{cases} \quad (4.3)$$

For convenience we set

$$\ell(u) = \int_\Omega g(x) u(x) dx \quad (4.4)$$

and for any $\mu > 0$ we denote by u_μ the weak solution to

$$\begin{cases} -\mu \Delta u_\mu + \beta_\infty u_\mu = f(x) \text{ in } \Omega, \\ u_\mu = 0 \text{ on } \partial\Omega. \end{cases} \quad (4.5)$$

As usual, solving a problem like (4.3) reduces to solve an equation in \mathbb{R} (see [4], [1]). Here arguing on $\ell(u)$ or $\alpha(\ell(u))$ offers two different equations. Indeed we have first

Theorem 4.1. *The mapping $u \rightarrow \ell(u)$ is a one-to-one mapping from the set of solutions to (4.3) into the set of solutions of the equation in \mathbb{R}*

$$\mu = \ell(u_{\alpha(\mu)}). \quad (4.6)$$

Proof. Suppose that u_∞ is solution to (4.3). Then, with our notation for u_μ

$$u_\infty = u_{\alpha(\ell(u_\infty))}$$

this implies

$$\ell(u_\infty) = \ell(u_{\alpha(\ell(u_\infty))})$$

i.e. $\ell(u_\infty)$ is solution to (4.6). Conversely, suppose that μ is solution to (4.6). Then, $u_{\alpha(\mu)}$ satisfies

$$\begin{cases} -\alpha(\mu) \Delta u_{\alpha(\mu)} + \beta_\infty u_{\alpha(\mu)} = f(x) \text{ in } \Omega, \\ u_{\alpha(\mu)} = 0 \text{ on } \partial\Omega. \end{cases}$$

Since, by (4.6), $\alpha(\mu) = \alpha(\ell(u_{\alpha(\mu)}))$, $u_{\alpha(\mu)}$ is solution to (4.3). The injectivity of the map $u \rightarrow \ell(u)$ is due to the fact that if $\ell(u_1) = \ell(u_2)$ when u_1 and u_2 are solutions to (4.3) then clearly $u_1 = u_2$. This completes the proof of the theorem. \square

It is now interesting to remark that the set of solutions can also be characterised by another set of fixed points namely :

Theorem 4.2. *The mapping $u \rightarrow \alpha(\ell(u))$ is a one-to-one mapping from the set of solutions to (4.3) into the set of solutions of the equation in \mathbb{R}*

$$\mu = \alpha(\ell(u_\mu)). \quad (4.7)$$

Proof. Suppose that u_∞ is solution to (4.3). Then, with our definition for u_μ

$$u_\infty = u_{\alpha(\ell(u_\infty))}$$

this implies that

$$\alpha(\ell(u_\infty)) = \alpha(\ell(u_{\alpha(\ell(u_\infty))}))$$

i.e. $\alpha(\ell(u_\infty))$ is solution to (4.7). Conversely, suppose that μ is solution to (4.7). Then u_μ is solution to

$$\begin{cases} -\alpha(\ell(u_\mu))\Delta u_\mu + \beta_\infty u_\mu = f(x) \text{ in } \Omega, \\ u_\mu = 0 \text{ on } \partial\Omega, \end{cases}$$

i.e. u_μ is solution to (4.3). To prove the injectivity of the map $u \rightarrow \alpha(\ell(u))$ one has just to notice that if $\alpha(\ell(u_1)) = \alpha(\ell(u_2))$ when u_1, u_2 are solutions to (4.3) then clearly $u_1 = u_2 = u_{\alpha(\ell(u_i))}$. This completes the proof of the theorem. \square

Then we can now show

Theorem 4.3. *Suppose that for some constants α_0, α_1 one has*

$$0 < \alpha_0 \leq \alpha \leq \alpha_1, \tag{4.8}$$

then the problem (4.3) admits at least one solution.

Proof. Due to (4.8) the strait line $y = \mu$ is cutting the curve $y = \alpha(\ell(\mu))$ and the result follows from the theorem 4.2. \square

Remark 2. *Of course (4.7) can have several solutions and even an infinity. In the case of a single solution it would be interesting and non trivial to show the convergence of $u(t)$ toward u_∞ . In the next paragraph we address a simple case to show what is on stake. We made it voluntary simple in a didactic spirit.*

Let us suppose that g is an eigenvalue of the Dirichlet problem i.e. that for some $\lambda > 0$, g satisfies in a weak sense

$$-\Delta g = \lambda g \text{ in } \Omega, \quad g = 0 \text{ on } \partial\Omega. \tag{4.9}$$

Then we have

Theorem 4.4. *Let g be solution to (4.9). Suppose that β is a positive constant, $(f, g) > 0$ and that the equation*

$$(\lambda\alpha(\mu) + \beta)\mu = (f, g) \tag{4.10}$$

admits a unique solution. Then if $u(x, t)$ is solution to (1.1) and u_∞ solution to (4.3) one has

$$|u(x, t) - u_\infty|_2 \rightarrow 0 \text{ when } t \rightarrow +\infty. \tag{4.11}$$

Proof. It is enough to show (see [2]) that $\ell(u(x, t)) \rightarrow \ell(u_\infty)$ when $t \rightarrow \infty$. Multiplying the equation (1.1) by g and integrating on Ω one gets

$$\frac{d}{dt}(u, g) + \alpha(\ell(u)) \int_{\Omega} \nabla u \nabla g dx + \beta(u, g) = (f, g)$$

i.e. using the definition of g and ℓ it comes

$$\frac{d}{dt}\ell(u) + \lambda\alpha(\ell(u))\ell(u) + \beta\ell(u) = (f, g).$$

Denote by μ_∞ the unique solution to (4.10). Since we assume $(f, g) > 0$ one has $\mu_\infty > 0$ and $(\lambda\alpha(\mu) + \beta)\mu < (f, g)$ for $\mu < \mu_\infty$. Indeed, for $\mu = 0$ $(\lambda\alpha(\mu) + \beta)\mu < (f, g)$ and the inequality follows for any $\mu < \mu_\infty$ since the solution to (4.10) is supposed to be unique. Suppose that

$$\ell(u_0) < \mu_\infty.$$

Since $\ell(u)$ is solution to the differential equation

$$\frac{d}{dt}\ell(u) = (f, g) - (\lambda\alpha(\ell(u)) + \beta)\ell(u)$$

$\ell(u)$ is increasing and of course converging toward μ_∞ . Similarly $\ell(u_0) > \mu_\infty$ implies that $\ell(u)$ is decreasing toward μ_∞ . This completes the proof of the theorem. \square

Remark 3. *In the case that we just considered one could describe the asymptotic behaviour of u using the same argument when the equation (4.10) admits different isolated solutions. We leave the proof to the reader.*

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